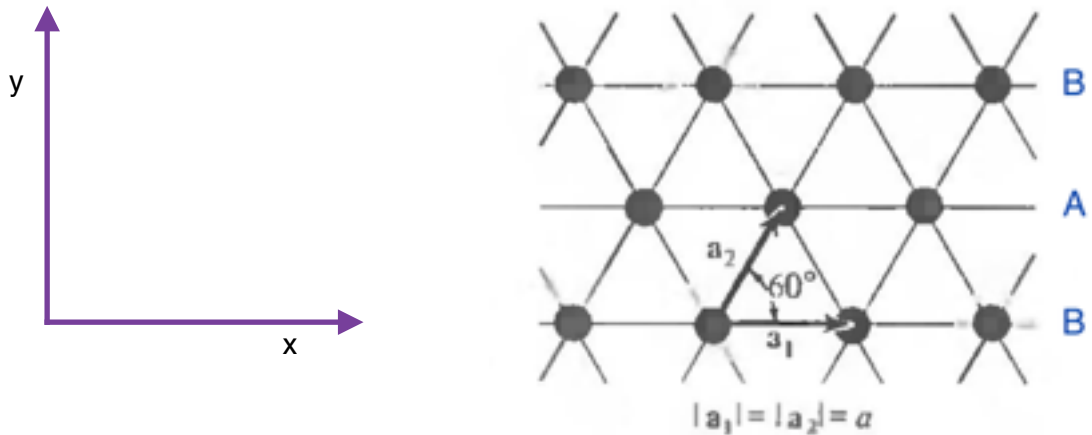


(time 3 hours)

Solve the following two exercises: (i) Give all details which help in understanding the proposed solution. Answers which only contain the final result or not enough detail will be judged insufficient and discarded; (ii) if you are requested to give evaluation/estimates, do so using 3 significant figures.

Exercise 1: Frustrated antiferromagnet: mean field

Consider N spins $1/2$, in the presence of a magnetic field $\mathbf{H} = H\hat{y}$, at the sites of a triangular lattice (2 dimensions!) with nearest neighbours antiferromagnetic coupling $J = -K_B T_N$ and K_B the Boltzmann constant and T_N a temperature. The triangular lattice can be seen as a collection of one dimensional lattices (say in the x direction) with lattice parameter a , parallel to each other. The lattice has area A , so that the density of sites is $n = N/A$. Moreover that lattice can be seen as made by two sublattices, A and B , made by lines of type A and B respectively.



1. First of all consider independent spins ($1/2$) in the external magnetic field $\mathbf{H} = H\hat{y}$ and compute the free energy of a spin in the presence of the field \mathbf{H} .
2. Using the result of the previous point, compute the average magnetization \mathcal{M} of a single spin.
3. Write down the hamiltonian of the interacting system.
4. Make a mean field approximations, assuming that the spins in the alternating horizontal *lines* A and B have average values $\langle S_x \rangle = 0$ and $(n\mu_B g/2)\langle S_y \rangle = M_A$ on sublattice A and $\langle S_x \rangle = 0$ and $(n\mu_B g/2)\langle S_y \rangle = M_B$ on sublattice B . Calculate the effective magnetic fields h_A and h_B , respectively at a site A and B . Here, M_A and M_B are sublattice magnetization densities: $M_A = (n/2) \mathcal{M}_A$, $M_B = (n/2) \mathcal{M}_B$, and \mathcal{M}_A and \mathcal{M}_B are single spin magnetizations, as calculated at point 2 above. NOTE: the effective field at a site is obtained adding to the external field the interaction with the 6 nearest neighbours, some of which are on sublattice A and some on sublattice B .
5. Write down the two selfconsistent equations for M_A and M_B when $H=0$. Remember that \mathcal{M}_A and \mathcal{M}_B are each function of the effective fields at the chosen sublattice.
6. Assuming that $M_A = -M_B$, show that there exists a critical temperature T_c , below which non zero sub lattice magnetizations are predicted.

Exercise 2: Lattice specific heat in D dimensions

Consider a lattice in dimension D , with *acoustic* phonon branches with small wave vector dispersion $\omega_s(\mathbf{k}) = \omega_{0s} (k/q_s)^\nu$, $s=1,2,.. D$, $\nu > 0$ and ω_0 and q_s a given frequency and a given wave vector. Denote with Ω_D the *solid angle* in D dimensions.

1. Express the energy density (energy per unit volume) due to phonons at finite temperature in terms of an integral containing the phonon density of states $g(\omega)$ and the Plank distribution.
2. Demonstrate that as $T \rightarrow 0$ the energy density is determined solely by the density of states for $\omega \rightarrow 0$.
3. As the optical branches (if there is more than one atom in the basis) have finite energy for every \mathbf{k} in the FBZ, at small energy (frequency) the phononic density of states $g(\omega)$ is completely determined by the acoustic branches at small \mathbf{k} . Calculate $g(\omega)$ at small frequencies.
4. Use the density of states found above in the expression for the energy density found at point 1 and extend the frequency integral to infinity, motivating the accuracy of such a choice.
5. Manipulate the expression for the energy density found at the previous point in such a way to obtain an integral that is temperature independent, and discuss its convergence properties.
6. Obtain the dependence on temperature, as $T \rightarrow 0$, of the lattice specific heat.